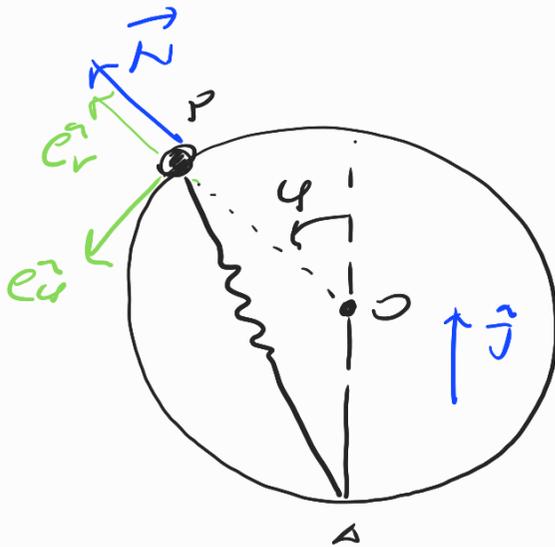


Ejercicio 7)



Dado que $\vec{N} \cdot \vec{v} = 0$, la energía de la partícula se conserva!

$$E = \frac{1}{2} m \vec{v}^2 + U_{el}, \quad \vec{v} = R \dot{\varphi} \vec{e}_{\varphi}, \quad U_{el} = \frac{1}{2} K (P-A)^2$$

$$P-A = P-O + O-A = R \vec{e}_r + R \hat{j}$$

$$\begin{aligned} (P-A)^2 &= 2R^2(\gamma + \hat{j} \cdot \vec{e}_r) = \\ &= 2R^2(\gamma + \cos\varphi) \end{aligned}$$

$$E = \frac{1}{2} m R^2 \dot{\varphi}^2 + \frac{1}{2} K 2R^2(\gamma + \cos\varphi)$$

$$E(\omega) = \frac{1}{2} m v \omega^2 + \frac{1}{2} K (2R)^2$$

$$\left| \frac{1}{2} m R^2 \dot{\varphi}^2 + K R^2(\gamma + \cos\varphi) = \frac{1}{2} m v \omega^2 + 2K R^2 \right|$$

b) 2da ley de Newton proyectada según \vec{e}_r

$$-m R \dot{\varphi}^2 = N - \underbrace{K(P-A) \cdot \vec{e}_r}_{R(\gamma + \hat{j} \cdot \vec{e}_r)}$$

$$N = K R (\gamma + \cos\varphi) - m R \dot{\varphi}^2 ;$$

Despejando $\dot{\varphi}^2$ de la ec. de momentos y substituyamos

en la anterior:

$$N(\varphi) = 3KR \cos \varphi - KR - \frac{mv_0^2}{R}$$

para $\varphi = \varphi_0$ $N(\varphi_0) = 0!$ $\left[\cos \varphi_0 = \frac{1}{3} + \frac{mv_0^2}{3KR^2} \right] \quad (7)$

si: $v_0^2 < \frac{2KR^2}{m}$

si $v_0^2 \geq \frac{2KR^2}{m}$ la partícula se desprende en el instante inicial



La partícula se mueve bajo la acción de una fuerza central conservativa $\left\{ \begin{array}{l} l = \text{cte.} \\ E = \text{cte.} \end{array} \right.$

$$l = m r^2 \dot{\varphi}$$

$$E = \frac{1}{2} m \vec{v}^2 + U_{el} = \frac{1}{2} m \dot{r}^2 + \frac{1}{2} m r^2 \dot{\varphi}^2 + \frac{1}{2} K r^2 \quad \left. \vphantom{E} \right\} \rightarrow$$

$$\rightarrow E = \frac{1}{2} m \dot{r}^2 + \frac{1}{2} m r^2 \left(\frac{l}{m r^2} \right)^2 + \frac{1}{2} K r^2 =$$

$$\frac{1}{2} m \dot{r}^2 + \frac{l^2}{2 m r^2} + \frac{1}{2} K r^2 :$$

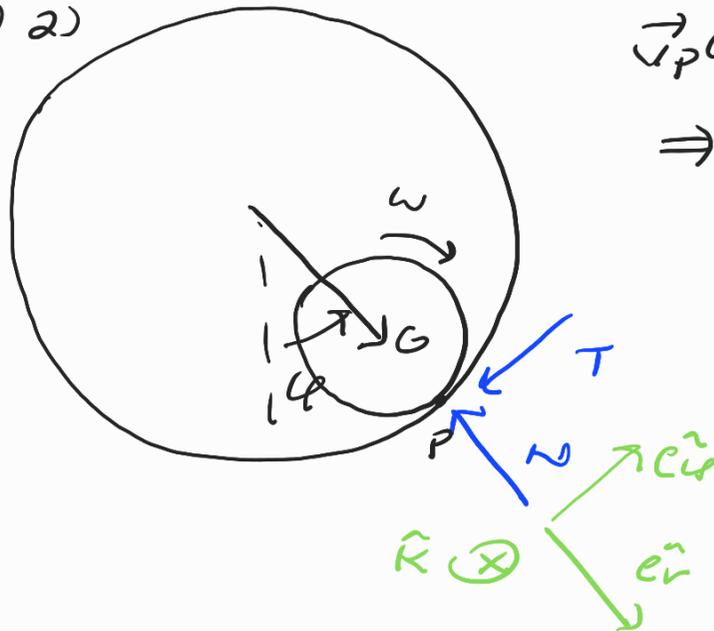
donde!

$$l = m 2 R v_0$$

$$E = \frac{1}{2} m v_0^2 + 2 K R^2$$

$$\dot{r}^2 = \frac{2}{m} \left[E - \frac{l^2}{2 m r^2} - \frac{1}{2} K r^2 \right]$$

Ejercicio 2) 2)



$$\vec{v}_P(\omega) = v_0 \hat{e}_\varphi$$

$$\Rightarrow \vec{T} = -T \hat{e}_\varphi,$$

$$T = f_0 N$$

1^{er} Condición al disco:

$$e_{\vec{r}}) - N = -m 3r \dot{\varphi}^2 \quad (i)$$

$$e_{\vec{\varphi}}) - T = m 3r \ddot{\varphi} \quad (ii)$$

usando $T = f_0 N$ y eliminando N entre las anteriores:

$$\boxed{\ddot{\varphi} + f_0 \dot{\varphi}^2 = 0} \quad (I)$$

2^{da} Condición sobre G :

$$I_G, \hat{K} \dot{\omega} = r T:$$

$$\frac{1}{2} m r^2$$

$$\frac{1}{2} m r \dot{\omega} = T \quad (iii)$$

Eliminando T entre (ii) y (iii):

$$\boxed{\dot{\omega} + 6\ddot{\varphi} = 0} \quad (II)$$

b. I) Integrando (II) en el tiempo:

$$\omega - \omega(\omega) + 6(\dot{\varphi} - \dot{\varphi}(\omega)) = 0$$

$$\Rightarrow \omega = -\alpha \dot{\varphi} + 2 \frac{v_0}{r}$$

$\left(\begin{array}{l} \omega(0) = 0 \\ \dot{\varphi}(0) = \frac{v_0}{3r} \end{array} \right)$

II) La ecuación (I) se puede escribir como:

$$\frac{d\dot{\varphi}}{dt} = -f_0 \dot{\varphi}^2 \rightarrow$$

integrando en variables

separadas:

$$\int_{\dot{\varphi}(0)}^{\dot{\varphi}} \frac{d\dot{\varphi}}{\dot{\varphi}^2} = -f_0 t :$$

$\frac{v_0}{3r}$

$$\frac{3r}{v_0} - \frac{1}{\dot{\varphi}} = -f_0 t : \dot{\varphi} = \left(f_0 t + \frac{3r}{v_0} \right)^{-1}$$

c) La velocidad del pto. P es:

$$\vec{v}_P = \vec{v}_G + \vec{\omega} \times (P-G) = (3r\dot{\varphi} - \omega r) \hat{e}_\varphi$$

$$= \left[3r\dot{\varphi} - \left(-\alpha\dot{\varphi} + 2\frac{v_0}{r} \right) r \right] \hat{e}_\varphi =$$

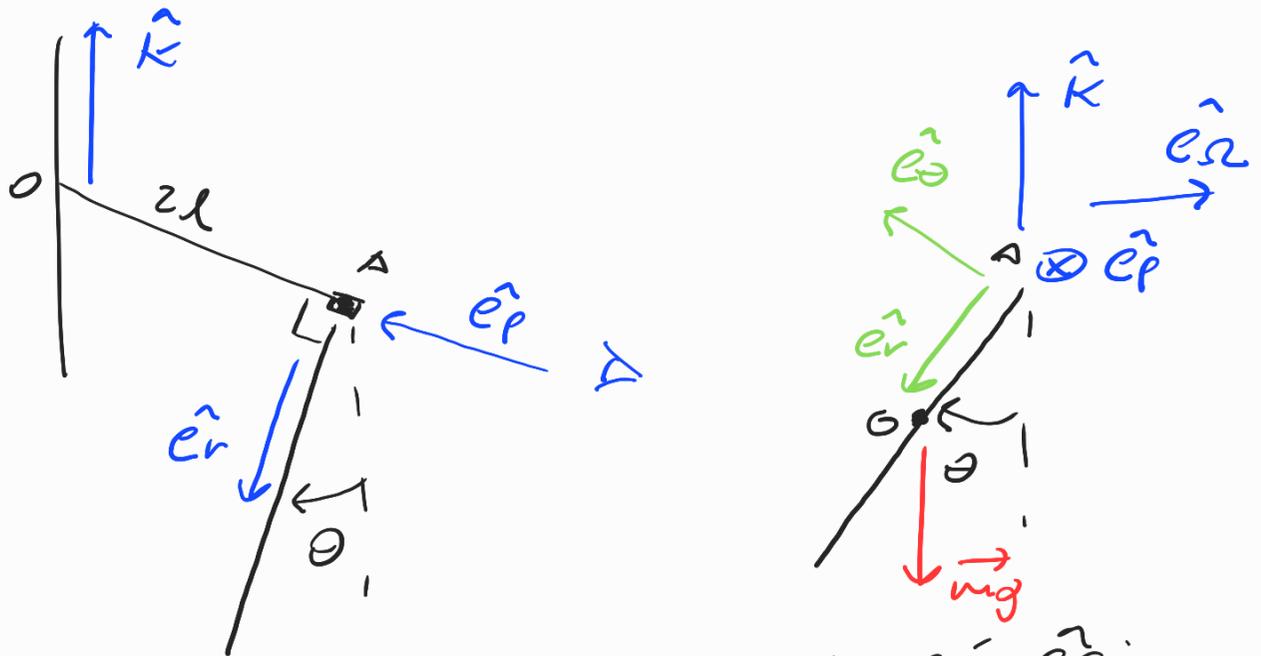
$$(9r\dot{\varphi} - 2v_0) \hat{e}_\varphi$$

Para $t=t_d$ el disco comienza a rodar sin deslizarle:

$$9r\dot{\varphi} - 2v_0 = 0 : \dot{\varphi}^{-1} = \frac{9}{2} \frac{r}{v_0} :$$

$$f_0 t_d + \frac{3r}{v_0} = \frac{9}{2} \frac{r}{v_0} : \boxed{t_d = \frac{3}{2} \frac{r}{v_0 f_0}}$$

Ejercicio 3)



2) 2^{da} Condición en A a la barra proyectada según \hat{e}_θ :

$$\left(m(G-A) \times \vec{a}_A + \frac{d}{dt} (\mathbb{I}_A \vec{\omega}) = \vec{M}_A^{(ext)} \right) \cdot \hat{e}_\theta,$$

$$\vec{M}_A^{(ext)} \cdot \hat{e}_\theta = \vec{M}_A^{(res)} \cdot \hat{e}_\theta \quad (\text{ya que } \vec{M}_A^{(rot)} \cdot \hat{e}_\theta = 0)$$

$$= -mgl \sin \theta;$$

$$\left. \begin{array}{l} G-A = l \hat{e}_r \\ \vec{a}_A = 2l \Omega^2 \hat{e}_\theta \end{array} \right\} m(G-A) \times \vec{a}_A \cdot \hat{e}_\theta = 0$$

$$\mathbb{I}_A \vec{\omega} : \mathbb{I}_A (\hat{e}_r, \hat{e}_\theta, \hat{e}_\phi) = \begin{pmatrix} 0 & 0 & 0 \\ 0 & I_{A, \hat{e}_\theta} & 0 \\ 0 & 0 & I_{A, \hat{e}_\phi} \end{pmatrix},$$

$$I_{A, \hat{e}_\theta} = I_{A, \hat{e}_\phi} = \int_0^{2l} \left(\frac{m}{2l} \right) dx x^2 = \frac{4}{3} ml^2;$$

$$\vec{\omega} = \Omega \hat{k} + \dot{\theta} \hat{e}_\rho = \Omega (-\cos\theta \hat{e}_r + \sin\theta \hat{e}_\theta) + \dot{\theta} \hat{e}_\rho$$

$$\Rightarrow \mathbb{I}_D \vec{\omega} = \Omega \left(\underbrace{-\cos\theta \mathbb{I}_D \hat{e}_r}_{=0} + \underbrace{\sin\theta \mathbb{I}_D \hat{e}_\theta}_{\frac{4}{3} m l^2 \hat{e}_\theta} \right) + \dot{\theta} \underbrace{\mathbb{I}_D \hat{e}_\rho}_{\frac{4}{3} m l^2 \hat{e}_\rho}$$

$$= \frac{4}{3} m l^2 \left[\Omega \sin\theta \hat{e}_\theta + \dot{\theta} \hat{e}_\rho \right];$$

$$\frac{d}{dt} (\mathbb{I}_D \vec{\omega}) \cdot \hat{e}_\rho = \frac{4}{3} m l^2 \left[\underbrace{\Omega \sin\theta \dot{\hat{e}}_\theta \cdot \hat{e}_\rho}_{\Omega \hat{k} \times \hat{e}_\theta \cdot \hat{e}_\rho = -\Omega \cos\theta} + \ddot{\theta} \right]$$

$$\Rightarrow \ddot{\theta} + \left(\frac{3}{4} g/l - \Omega^2 \cos\theta \right) \sin\theta = 0 \quad \rightarrow \frac{dV_{\text{eff}}}{d\theta}$$

b) Las posiciones de equilibrio corresponden a $\ddot{\theta} = 0 \quad \left(\frac{dV_{\text{eff}}}{d\theta} = 0 \right)$

$$\sin\theta \left[\frac{3}{4} g/l - \Omega^2 \cos\theta \right] = 0 :$$

$$\begin{cases} \sin\theta = 0 : \theta = 0, \pi \\ \cos\theta = \frac{\frac{3}{4} g/l}{\Omega^2} \end{cases}$$

c) $\theta_{eq} = 0, \pi$ existen siempre

$$\cos\theta_{eq} = \left(\frac{3}{4} g/l \right) / \Omega^2 \leq 1 : \exists \text{ si } \Omega^2 \geq \Omega_c^2 = \frac{3}{4} g/l$$

$$\text{Estabilidad : } \frac{d^2 V_{\text{eff}}}{d\theta^2} = \Omega^2 + (\Omega_c^2 - 2\Omega^2 \cos\theta) \cos\theta =$$

$$\begin{cases} (\theta_{eq} = 0) \quad \Omega_c^2 - \Omega^2 \geq 0 \text{ si } \Omega^2 \leq \Omega_c^2 : \text{estable si } \Omega^2 \leq \Omega_c^2 \\ (\theta_{eq} = \pi) \quad -(\Omega_c^2 + \Omega^2) < 0 : \text{siempre inestable} \\ (\cos\theta_{eq} = \frac{\Omega_c^2}{\Omega^2}) \quad \frac{\Omega_c^4 - \Omega_c^4}{\Omega^2} \geq 0 \text{ si } \Omega^2 \geq \Omega_c^2 \text{ estable inestable} \\ \text{exista} \end{cases}$$