

$$\hat{k} = -\cos\theta\hat{e}_1 + \sin\theta\hat{e}_\theta$$

$$\hat{z} = \sin\theta\hat{e}_1 + \cos\theta\hat{e}_\theta$$

a)  $\vec{L}_O^T = \vec{L}_O^A + \vec{L}_O^m$

masa  $\vec{r} = R\hat{z} + R\sin\theta\hat{e}_1$

$$\vec{v}_p = R\dot{\phi}\hat{j} + R(\dot{\phi}\sin\theta\hat{j} + \dot{\theta}\hat{e}_\theta)$$

$$\vec{L}_O^m = \vec{r} \wedge m\vec{v}_p$$

$$= (R\hat{z} + R\sin\theta\hat{e}_1) \wedge m [R\dot{\phi}(1+\sin\theta)\hat{j} + R\dot{\theta}\hat{e}_\theta]$$

$$= mR^2 (\dot{\phi}(1+\sin\theta)\hat{k} + mR^2\dot{\theta}\sin\theta(-\hat{j}) + mR^2\dot{\phi}(1+\sin\theta)\hat{e}_\theta - mR^2\dot{\theta}\hat{j})$$

$$= mR^2 [\dot{\phi}(1+\sin\theta)\hat{k} - \dot{\theta}(1+\sin\theta)\hat{j} + \dot{\phi}(1+\sin\theta)\hat{e}_\theta]$$

ARO  $\mathbb{I}_C^{\{\hat{e}_1, \hat{j}, \hat{k}\}} = m_a R^2 \begin{pmatrix} \frac{1}{2} & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & \frac{1}{2} \end{pmatrix}$

por completitud se agrega esta parte pero  $m_a \rightarrow 0$

$$\mathbb{I}_O = \mathbb{I}_C + \mathbb{I}_p^{(m,c)}$$

$$\vec{r}_c - \vec{r}_O = R\hat{z}$$

$$\mathbb{I}_p^{m,a} = m_a R^2 \begin{pmatrix} 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$\mathbb{I}_O = m_a R^2 \begin{pmatrix} \frac{1}{2} & 0 & 0 \\ 0 & 2 & 0 \\ 0 & 0 & \frac{3}{2} \end{pmatrix}$$

$$\vec{L}_O^A = \mathbb{I}_O \vec{\omega}_A \quad \vec{\omega}_A = \dot{\phi}\hat{k} \Rightarrow \vec{L}_O^A = \frac{3}{2} m_a R^2 \dot{\phi}\hat{k}$$

$$\Rightarrow \vec{L}_0 = R^2 \left[ \left( m(1+\sin\theta)\dot{\varphi} + \frac{3}{2} m_a \dot{\varphi} \right) \hat{k} - m\dot{\theta}(1+\sin\theta)\hat{j} + m\dot{\varphi}(1+\sin\theta)\hat{e}_\theta \right]$$

$$b) \vec{M}_0^{ext} = R\hat{e}_\theta (-mg\hat{k}) + (R\hat{e}_\theta + R\hat{e}_\theta) \wedge (-mg\hat{k}) + M_1^{react} \hat{e}_\theta + M_2^{react} \hat{j}$$

Ya que las Fuerzas entre la masa y el Aro forman un par de acción y reacción y se cancelan mutuamente.

De otra forma, son internas al sistema. total. giro libre según  $\hat{k}$

$$\vec{M}_0^{ext} = mgR\hat{j} + mgR\hat{j} + mgR\sin\theta\hat{j} + \overbrace{M_1^{react}\hat{e}_\theta + M_2^{react}\hat{j}}$$

$$\Rightarrow \vec{M}_0^{ext} \cdot \hat{k} = 0 \quad \rightarrow \quad \frac{d\vec{L}_0}{dt} \cdot \hat{k} = \vec{M}_0^{ext} \cdot \hat{k} = 0$$

$$\left. \begin{aligned} \frac{d(\vec{L}_0 \cdot \hat{k})}{dt} &= \frac{d\vec{L}_0}{dt} \cdot \hat{k} + \vec{L}_0 \cdot \frac{d\hat{k}}{dt} \\ &= \frac{d\vec{L}_0}{dt} \cdot \hat{k} \quad \left[ \begin{array}{l} \hat{k} \text{ es fijo} \\ \frac{d\hat{k}}{dt} = 0 \end{array} \right] \end{aligned} \right\} \rightarrow \frac{d(\vec{L}_0 \cdot \hat{k})}{dt} = 0$$

$$\Rightarrow \vec{L}_0 \cdot \hat{k} = L_z \text{ cte.}$$

$$L_z = R^2 \left( m(1+\sin\theta)\dot{\varphi} + \frac{3}{2} m_a \dot{\varphi} \right) + R^2 m \dot{\varphi} (1+\sin\theta) \overbrace{\hat{e}_\theta \cdot \hat{k}}^{\sin\theta}$$

$$L_z = R^2 \left( m \dot{\varphi} (1+\sin\theta)^2 + \frac{3}{2} m_a \dot{\varphi} \right)$$

Ahora  $m_a \rightarrow 0 \Rightarrow L_z = m R^2 \dot{\varphi} (1+\sin\theta)^2$

d) Fuerzas conservativas: Peso

$$\frac{dE}{dt} = P^{res} \quad P^{res} = P_{Normal}$$

$$\Rightarrow P_{Normal} = \vec{N}_A \cdot \vec{v}_A + \vec{N}_P \cdot \vec{v}_P$$

Hay fuerzas Normales actuando sobre el Aro ( $\vec{N}_A$ ) y sobre la partícula ( $\vec{N}_P$ ) por acción y reacción.  $\vec{N}_A = -\vec{N}_P$

$$\vec{v}_P = \vec{v}' + \vec{v}_T \text{ (Roverbol)} \quad \text{y} \quad \vec{v}_T = \vec{v}_A$$

$$\rightarrow P_{Normal} = \vec{N}_A \cdot (\vec{v}_A - \vec{v}' - \vec{v}_A) = \vec{N}_A \cdot \vec{v}' = 0 \quad \rightarrow P^{res} = 0 \text{ E cte}$$

0 pues  $m_a \rightarrow 0$

$$T = \frac{1}{2} m \vec{v}_P^2 + \frac{1}{2} \vec{\omega}_A \cdot \mathbb{I}_O \vec{\omega}_A \quad \rightarrow \quad T = \frac{1}{2} m \vec{v}_P^2$$

$$U_g = -mgR \cos\theta$$

$$\vec{v}_P = R(1+\sin\theta)\hat{j} + R\dot{\theta}\hat{e}_\theta$$

$$T = \frac{1}{2} m R^2 (1+\sin\theta)^2 \dot{\varphi}^2 + \frac{1}{2} m R^2 \dot{\theta}^2$$

$$E = E_0 = \frac{1}{2} m R^2 \left( \dot{\theta}^2 + \dot{\varphi}^2 (1 + \sin^2 \theta) \right) - mg R \cos \theta$$

$$E_0 = \frac{1}{2} m R^2 \dot{\varphi}_0^2 - mg R$$

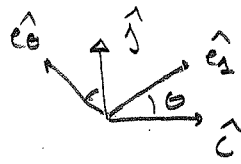
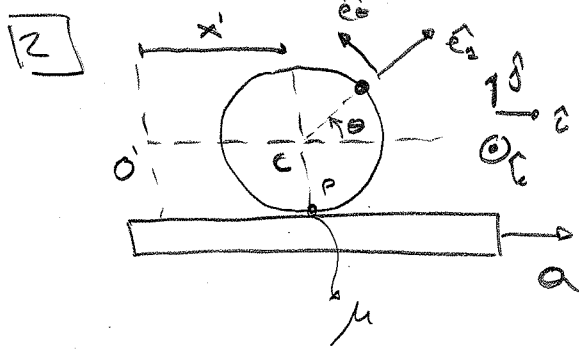
de la conservación de  $L_z$ :

$$L_z = m R^2 \dot{\varphi} (1 + \sin^2 \theta) \quad \rightarrow \quad L_z = m R^2 \dot{\varphi}_0$$

$$\rightarrow \dot{\varphi} = \frac{\dot{\varphi}_0}{(1 + \sin^2 \theta)^2}$$

$$c) \quad E_0 = \frac{1}{2} m R^2 \left( \dot{\theta}^2 + \frac{\dot{\varphi}_0^2}{(1 + \sin^2 \theta)^2} (1 + \sin^2 \theta)^2 \right) - mg R \cos \theta = \frac{1}{2} m R^2 \dot{\varphi}_0^2 - mg R$$

$$\dot{\theta}^2 = \dot{\varphi}_0^2 \left( 1 - \frac{1}{(1 + \sin^2 \theta)^2} \right) + \frac{2g}{R} (\cos \theta - 1)$$



$$\hat{e}_1 = \cos\theta \hat{c} + \sin\theta \hat{j}$$

$$\hat{e}_2 = -\sin\theta \hat{c} + \cos\theta \hat{j}$$

a) El centro de masa del rígido coincide con la partícula (el disco tiene masa despreciable)

$$\vec{r}_G = \vec{r}_C + R \hat{e}_1$$

$$\dot{\hat{e}}_1 = \bar{\omega} \wedge \hat{e}_1 = \dot{\theta} \hat{k} \wedge \hat{e}_1 = \dot{\theta} \hat{e}_2$$

$$\text{con } \vec{r}_C = \vec{r}_{O'} + x' \hat{c}$$

$$\dot{\hat{e}}_2 = \bar{\omega} \wedge \hat{e}_2 = \dot{\theta} \hat{k} \wedge \hat{e}_2 = -\dot{\theta} \hat{e}_1$$

$$\vec{v}_G = \vec{v}_C + R \dot{\hat{e}}_1 = \vec{v}_{O'} + \dot{x}' \hat{c} + R \dot{\theta} \hat{e}_2$$

$$\vec{a}_G = \vec{a}_{O'} + \ddot{x}' \hat{c} + R \ddot{\theta} \hat{e}_2 + R \dot{\theta} \dot{\hat{e}}_2 \quad \text{O' es solidario a la placa} \Rightarrow \vec{a}_{O'} = a \hat{c}$$

$$\vec{a}_G = a \hat{c} + \ddot{x}' \hat{c} + R \ddot{\theta} \hat{e}_2 - R \dot{\theta}^2 \hat{e}_1$$

RSD  $\vec{v}_P = 0$  (en el sistema solidario a la placa)

Dist

$$\text{velocidades } \vec{v}_P = \vec{v}_C + \bar{\omega} \wedge (\vec{r}_P - \vec{r}_C) \quad \vec{r}_P - \vec{r}_C = -R \hat{j}$$

$$\vec{v}_P = \dot{x}' \hat{c} + \dot{\theta} \hat{k} \wedge (-R \hat{j}) = (\dot{x}' + R \dot{\theta}) \hat{c} = 0 \Rightarrow \dot{x}' = -R \dot{\theta}$$

$$\downarrow$$

$$\ddot{x}' = -R \ddot{\theta}$$

En el sistema absoluto,  $\vec{v}_P = \vec{v}_{O'}$  punto de la placa

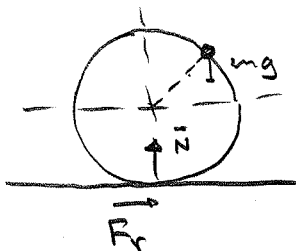
Dist velocidades

$$\vec{v}_P = \vec{v}_C + \dot{\theta} \hat{k} \wedge (-R \hat{j}) \quad \vec{v}_C = \vec{v}_{O'} + \dot{x}' \hat{c}$$

$$\Rightarrow \vec{v}_P = \vec{v}_{O'} + \dot{x}' \hat{c} + R \dot{\theta} \hat{c} = \vec{v}_{O'} \Rightarrow \dot{x}' = -R \dot{\theta} \text{ mismo resultado}$$

$$\vec{a}_G = a \hat{c} - R \ddot{\theta} \hat{c} + R \ddot{\theta} \hat{e}_2 - R \dot{\theta}^2 \hat{e}_1$$

b)



$$1^{\text{ra}} \text{ Cardinal } m \vec{a}_G = N \hat{j} + F_r \hat{c} - mg \hat{j}$$

$$\vec{a}_G = a\hat{e} - R\ddot{\theta}\hat{e} + R\dot{\theta}(-\sin\theta\hat{e} + \cos\theta\hat{j}) - R\dot{\theta}^2(\cos\theta\hat{e} + \sin\theta\hat{j})$$

1<sup>ra</sup> cardinal i)  $m(a - R\ddot{\theta}(1 + \sin\theta) - R\dot{\theta}^2\cos\theta) = F_r$  (1)

j)  $mR(\ddot{\theta}\cos\theta - \dot{\theta}^2\sin\theta) = N - mg$  (2)

2<sup>da</sup> cardinal en C (C tiene aceleración)

$$\vec{L}_C = m(\vec{r}_G - \vec{r}_C) \wedge \vec{v}_C + \mathbb{I}_C \vec{\omega}$$

$$\vec{r}_G - \vec{r}_C = R\hat{e}_1 = R\cos\theta\hat{e} + R\sin\theta\hat{j}$$

$$\vec{v}_C = \vec{v}_0 + \dot{x}'\hat{e} \quad \vec{a}_C = \vec{a}_0 + \ddot{x}'\hat{e} = a\hat{e} + \ddot{x}'\hat{e}$$

$$\frac{d\vec{L}_C}{dt} = m(\vec{v}_G - \vec{v}_C) \wedge \vec{v}_C + m(\vec{r}_G - \vec{r}_C) \wedge \vec{a}_C + \mathbb{I}_C \dot{\vec{\omega}}$$

La 2<sup>da</sup> cardinal:  $\frac{d\vec{L}_C}{dt} = m\vec{v}_G \wedge \vec{v}_C + \vec{M}_C^{(ext)}$

Juntando ambas:  $m(\vec{r}_G - \vec{r}_C) \wedge \vec{a}_C + \mathbb{I}_C \dot{\vec{\omega}} = \vec{M}_C^{ext}$

$$m(\vec{r}_G - \vec{r}_C) \wedge \vec{a}_C = -mR\sin\theta(a + \ddot{x}')\hat{k}$$

$$\mathbb{I}_C \{\hat{e}_1, \hat{e}_2, \hat{k}\} = mR^2 \begin{pmatrix} 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \quad \vec{\omega} = \dot{\theta}\hat{k} \quad \dot{\vec{\omega}} = \ddot{\theta}\hat{k}$$

$$\vec{M}_C^{ext} = (mR^2\ddot{\theta} - mR(a + \ddot{x}')\sin\theta)\hat{k}$$

de la dinámica

$$\vec{M}_C^{ext} = -mgR\cos\theta\hat{k} + F_r R\hat{k}$$

$$\Rightarrow mR^2\ddot{\theta} - mR(a + \ddot{x}')\sin\theta = RF_r - mgR\cos\theta \quad (3)$$

Sustituyendo  $F_r$  de (1) en (3)

$$mR(R\ddot{\theta} - (a + \ddot{x}')\sin\theta) = mR(a - R\ddot{\theta}(1 + \sin\theta) - R\dot{\theta}^2\cos\theta - g\cos\theta)$$

sustituyo  $\ddot{x}' = -R\ddot{\theta}$

Ecuación de movimiento.

c) En eq. relativo  $\dot{x}' = 0$   $\dot{\theta}' = 0$  y  $\ddot{x}' = 0$   $\ddot{\theta} = 0$

en la ec. de movimiento.

$$-mRa\sin\theta = mR(a - g\cos\theta) \rightarrow -a\sin\theta - a = -g\cos\theta$$

$$\rightarrow a(1 + \sin\theta) = g\cos\theta \quad \text{en } \theta = \frac{\pi}{4} \quad \sin\theta = \cos\theta = \frac{1}{\sqrt{2}}$$

$$\Rightarrow \frac{a}{g} = \frac{\cos\theta}{1 + \sin\theta}$$

$$\rightarrow \boxed{\frac{a}{g} = \frac{1}{\sqrt{2} + 1}}$$

Imponiendo el equilibrio en (1) y (2)

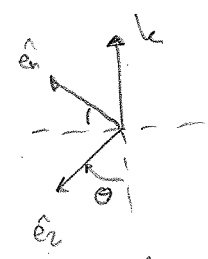
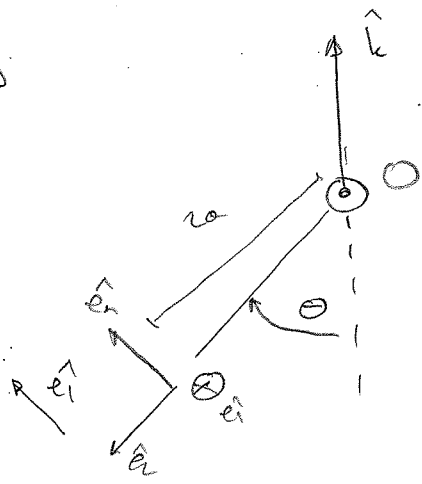
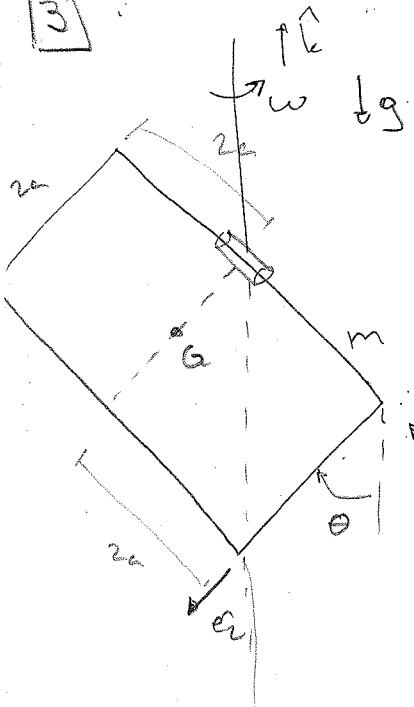
$$F_r = ma$$

$$N = mg$$

No desliza si  $|\vec{F}_r| \leq \mu |\vec{N}|$

$$\rightarrow a \leq \mu g \rightarrow \boxed{\frac{a}{g} \leq \mu}$$

3



$$\hat{k} = -\cos\theta \hat{e}_1 + \sin\theta \hat{e}_2$$

$$\vec{\omega} = \omega \hat{k} + \dot{\theta} \hat{e}_1$$

$$= -\omega \cos\theta \hat{e}_1 + \omega \sin\theta \hat{e}_2 + \dot{\theta} \hat{e}_1$$

a)  $\vec{L}_O = \Pi_O \vec{\omega}$

$$\Pi_O = \Pi_G + J_O^{m,G}$$

$$\vec{r}_G - \vec{r}_O = a \hat{e}_2$$

$$\Pi_G^{(\hat{e}_1, \hat{e}_2, \hat{e}_3)} = \frac{m}{12} \begin{pmatrix} (2a)^2 & 0 & 0 \\ 0 & (4a)^2 & 0 \\ 0 & 0 & (4a)^2 + (2a)^2 \end{pmatrix}$$

$$J_O^{m,G} = ma^2 \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$= \frac{m a^2}{3} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 4 & 0 \\ 0 & 0 & 5 \end{pmatrix}$$

$$\Pi_O = ma^2 \begin{pmatrix} 5/3 & 0 & 0 \\ 0 & 4/3 & 0 \\ 0 & 0 & 2/3 \end{pmatrix}$$

$$\vec{L}_O = \frac{4ma^2}{3} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 2 \end{pmatrix} \begin{pmatrix} \dot{\theta} \\ -\omega \cos\theta \\ \omega \sin\theta \end{pmatrix}$$

$$\vec{L}_O = \frac{4}{3} ma^2 (-\omega \cos\theta \hat{e}_1 + \omega \sin\theta \hat{e}_2 + 2\dot{\theta} \hat{e}_3)$$

$$+ \frac{4}{3} ma^2 \dot{\theta} \hat{e}_1$$

$$b) \dot{\vec{L}}_0 = \vec{M}_0^{ext}$$

$$\dot{\vec{L}}_0 = \frac{4}{3} m a^2 \left( + \omega \dot{\theta} \sin \theta \hat{e}_2 - \omega \cos \theta \hat{e}_2 + 2 \omega \dot{\theta} \cos \theta \hat{e}_n + 2 \omega \sin \theta \hat{e}_n \right. \\ \left. + \ddot{\theta} \hat{e}_1 + \dot{\theta} \hat{e}_1 \right)$$

$$\dot{\hat{e}}_1 = \vec{\omega} \wedge \hat{e}_1 = (-\omega \cos \theta \hat{e}_2 + \omega \sin \theta \hat{e}_n + \dot{\theta} \hat{e}_1) \wedge \hat{e}_1 = + \omega \cos \theta \hat{e}_n + \omega \sin \theta \hat{e}_2$$

$$\dot{\hat{e}}_2 = (-\omega \cos \theta \hat{e}_2 + \omega \sin \theta \hat{e}_n + \dot{\theta} \hat{e}_1) \wedge \hat{e}_2 = -\omega \sin \theta \hat{e}_1 + \dot{\theta} \hat{e}_n$$

$$\dot{\hat{e}}_n = (-\omega \cos \theta \hat{e}_2 + \omega \sin \theta \hat{e}_n + \dot{\theta} \hat{e}_1) \wedge \hat{e}_n = -\omega \cos \theta \hat{e}_1 - \dot{\theta} \hat{e}_2$$

$$\dot{\vec{L}}_0 = \frac{4}{3} m a^2 \left( \omega \dot{\theta} \sin \theta \hat{e}_2 + \omega^2 \cos \theta \sin \theta \hat{e}_1 + 2 \omega \dot{\theta} \cos \theta \hat{e}_n \right. \\ \left. + 2 \omega \sin \theta (-\omega \cos \theta) \hat{e}_1 + \dot{\theta} \hat{e}_1 + \dot{\theta} (\omega \cos \theta \hat{e}_n + \omega \sin \theta \hat{e}_2) \right. \\ \left. - \omega \dot{\theta} \cos \theta \hat{e}_n - 2 \omega \dot{\theta} \sin \theta \hat{e}_2 \right)$$

$$\vec{M}_0^{ext} = -mga \sin \theta \hat{e}_1 + M_2 \hat{e}_2 + M_n \hat{e}_n$$



No hay momentos reactivos según  $\hat{e}_1$ .

$$E_c \text{ de mov.: } \dot{\vec{L}}_0 \cdot \hat{e}_1 = \vec{M}_0^{ext} \cdot \hat{e}_1$$

$$\frac{4}{3} m a^2 \left( \ddot{\theta} - \omega^2 \sin \theta \cos \theta \right) = -mga \sin \theta$$

$$\ddot{\theta} - \omega^2 \sin \theta \cos \theta + \frac{3}{4} \frac{g}{a} \sin \theta = 0 \text{ es del tipo } \ddot{\theta} + f(\theta) = 0$$

puntos de eq  $\theta_{eq}$  t.q  $f(\theta_{eq}) = 0$  si  $\left. \frac{df}{d\theta} \right|_{\theta_{eq}} > 0$  estables

$\left. \frac{df}{d\theta} \right|_{\theta_{eq}} < 0$  inestables.

$$f(\theta) = \sin \theta \left( \frac{3}{4} \frac{g}{a} - \omega^2 \cos \theta \right) \rightarrow \theta_1 = 0, \\ \theta_2 = \pi$$

$$\theta_3 \text{ t.q } \cos \theta_3 = \frac{3}{4} \frac{g}{a} \frac{1}{\omega^2}$$

$$\theta_3 \exists \text{ ssi } \frac{3}{4} \frac{g}{a} \frac{1}{\omega^2} < 1$$



$$\frac{df}{d\theta} = \cos\theta \left( \frac{3g}{4a} - \omega^2 \cos\theta \right) + \omega^2 \sin^2\theta$$

$$\left. \frac{df}{d\theta} \right|_{\theta_1=0} = \frac{3g}{4a} - \omega^2 > 0 \text{ si } \frac{3g}{4a} \frac{1}{\omega^2} > 1 \text{ estable}$$

$$\text{si } \frac{3g}{4a} \frac{1}{\omega^2} < 1 \text{ inestable.}$$

$$\left. \frac{df}{d\theta} \right|_{\theta_2=\pi} = -1 \left( \frac{3g}{4a} + \omega^2 \right) < 0 \text{ siempre inestable siempre}$$

$$\left. \frac{df}{d\theta} \right|_{\theta_3} = \omega^2 \sin^2\theta_3 > 0 \text{ Estable}$$

$\Rightarrow \theta_3$  cuando Existe es estable.