

$$\vec{v}_T = \vec{v}_A + \vec{\omega}_2 \wedge \vec{r}'$$

$$\vec{\omega}_2 = \Omega \vec{k}$$

$$\vec{\omega}_2 \wedge \vec{r}' = \Omega \vec{k} \wedge (L \vec{e}_1 + x \vec{e}_3) = L \Omega \vec{e}_2 + x \Omega \vec{e}_4 = \vec{v}_T$$

$$\vec{v} = \dot{x} \vec{e}_3 + x(\dot{\theta} + \Omega) \vec{e}_4 + L \Omega \vec{e}_2 \quad \text{Mismo resultado de antes}$$

$$\vec{a}' = \ddot{x} \vec{e}_3 + \dot{x} \frac{d \vec{e}_3}{dt} + \dot{x} \ddot{\theta} \vec{e}_4 + x \ddot{\theta} \vec{e}_4 + x \dot{\theta} \frac{d \vec{e}_4}{dt} =$$

$$\ddot{\theta} \vec{e}_4 - \dot{\theta}^2 \vec{e}_3$$

$$\vec{a}' = (\ddot{x} - x \dot{\theta}^2) \vec{e}_3 + (2 \dot{x} \dot{\theta} + x \ddot{\theta}) \vec{e}_4$$

$$\vec{a}_G = 2 \vec{\omega}_2 \wedge \vec{v}' = 2 \Omega \vec{k} \wedge (\dot{x} \vec{e}_3 + x \dot{\theta} \vec{e}_4) = 2 \Omega \dot{x} \vec{e}_4 - 2 \Omega x \dot{\theta} \vec{e}_3 = \vec{a}_G$$

$$\vec{a}_T = \vec{a}_A + \vec{\omega}_2 \wedge \vec{r}' + \vec{\omega}_2 \wedge [\vec{\omega}_2 \wedge \vec{r}'] = \Omega^2 \vec{k} [\vec{k} \wedge (L \vec{e}_1 + x \vec{e}_3)] =$$

$$= \Omega^2 \vec{k} [L \vec{e}_2 + x \vec{e}_4] = \Omega^2 (-L \vec{e}_1 - x \vec{e}_3)$$

$$\vec{a}_T = -\Omega^2 (L \vec{e}_1 + x \vec{e}_3)$$

$$\vec{a} = \underbrace{(\ddot{x} - x \dot{\theta}^2 - 2 \Omega x \dot{\theta} - x \Omega^2)}_{-x(\dot{\theta} + \Omega)^2} \vec{e}_3 + [2 \dot{x}(\dot{\theta} + \Omega) + x \ddot{\theta}] \vec{e}_4 - L \Omega^2 \vec{e}_1$$

Mismo resultado de antes

Otro sistema relativo: $\{B, \vec{e}_1, \vec{e}_2, \vec{k}\}$

$$\vec{r}' = x \vec{e}_3 \Rightarrow \vec{v}' = \dot{x} \vec{e}_3 + x \frac{d \vec{e}_3}{dt} = \dot{x} \vec{e}_3 + x \dot{\theta} \vec{e}_4 = \vec{v}'$$

$$\dot{\theta} \vec{e}_4$$

$$\vec{v}_T = \vec{v}_B + \vec{\omega}_3 \wedge \vec{r}' = L \Omega \vec{e}_2 + x \Omega \vec{e}_4 = \vec{v}_T$$

$$L \Omega \vec{e}_2 \quad \Omega \vec{k}$$

Mismos \vec{v}' y \vec{v}_T como anteriores
 $\Rightarrow \vec{v}$ será el mismo

$$\vec{a}' = (\ddot{x} - x \dot{\theta}^2) \vec{e}_3 + (2 \dot{x} \dot{\theta} + x \ddot{\theta}) \vec{e}_4 \quad \text{Misma } \vec{a}' \text{ de antes}$$

$$\vec{a}_G = 2 \vec{\omega}_3 \wedge \vec{v}' = 2 \Omega \dot{x} \vec{e}_4 - 2 \Omega x \dot{\theta} \vec{e}_3 = \vec{a}_G$$

xg' es mismo $\vec{\omega}$ y \vec{v}' de caso anterior

$$\vec{a}_T = \vec{a}_B + \vec{\omega}_3 \wedge \vec{r}' + \vec{\omega}_3 \wedge [\vec{\omega}_3 \wedge \vec{r}']$$

$$-L \Omega^2 \vec{e}_1$$

$$x \Omega^2 \vec{k} \wedge (\vec{k} \wedge \vec{e}_3) = -x \Omega^2 \vec{e}_3$$

$$\vec{e}_4$$

según \vec{e}_3 : $m(\ddot{x} - x\Omega^2) = -Kx \Rightarrow \ddot{x} + x\Omega^2 = 0$

" \vec{e}_4 : $m(L\Omega^2 + 2\dot{x}\Omega) = N$

$x(t) = A\cos\Omega t + B\sin\Omega t$

$x(0) = L = A$

$\dot{x}(t) = -L\Omega\sin\Omega t + B\Omega\cos\Omega t$

$\dot{x}(0) = 0 \Rightarrow B = 0 \Rightarrow x(t) = L\cos\Omega t$

$\dot{x}(t) = -L\Omega\sin\Omega t$

$N(t) = mL\Omega^2(1 - 2\sin\Omega t) \geq 0$

$\sin\Omega t \leq \frac{1}{2}$

$\Omega t \leq \frac{\pi}{6} \Rightarrow t_{\text{dep}} = \frac{\pi}{6\Omega}$

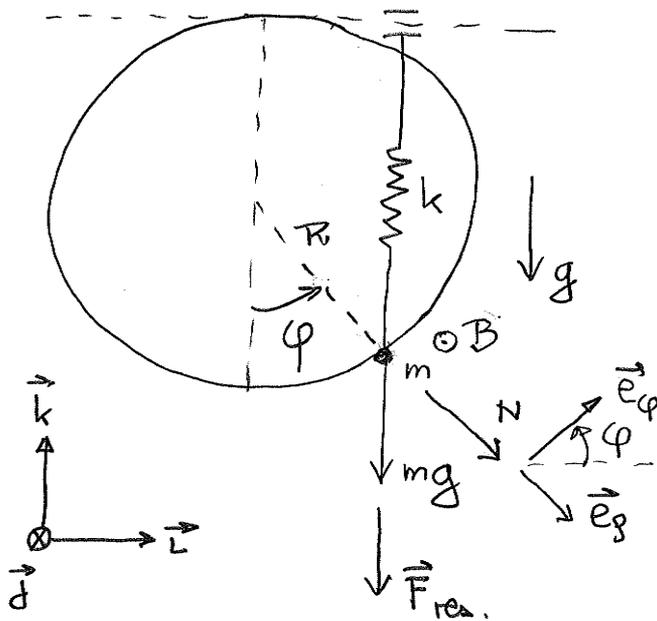
$x(t_{\text{dep}}) = L\cos\frac{\pi}{6} = \frac{\sqrt{3}}{2}L < 2L \Rightarrow$ la partícula no se escapa de la barra AB porque $x(t)$ disminuye.

Y la partícula tampoco llega al extremo B:

$x(t) = 0 \Rightarrow \cos\Omega t = 0 \Rightarrow \Omega t = \frac{\pi}{2}$

$t_B = \frac{\pi}{2\Omega} > t_{\text{dep}}$

Ejercicio N° 2



Guía lisa $\Rightarrow \vec{T} = 0$

$$\vec{F}_{rea} = -k(l - l_0)(-\vec{k}) = k l \vec{k}$$

$$l = R + R \cos \varphi$$

Vínculo Bilateral: $N \geq 0$

$$kR = mg$$

parte a:

1) Reacción de la guía lisa:

$$\vec{R} = N \vec{e}_\varphi - B \vec{j}$$

$$\vec{F} = R \vec{e}_\varphi$$

$$\vec{v} = R \dot{\varphi} \vec{e}_\varphi$$

$$\vec{R} \cdot \vec{v} = 0 \Rightarrow \text{Reacción de Potencia Nula}$$

2) Peso \Rightarrow conservativo: $U_g = mgz$ $z = \vec{r} \cdot \vec{k} = R \vec{e}_\varphi \cdot \vec{k} = -R \cos \varphi$

$$\Rightarrow U_g = -mgR \cos \varphi$$

3) Resorte \Rightarrow conservativo: $U_k = \frac{k l^2}{2} = \frac{k (R + R \cos \varphi)^2}{2}$

$$U_k = \frac{k R^2 (1 + \cos \varphi)^2}{2}$$

No hay fuerzas residuales \Rightarrow sistema conservativo.

parte b: Sist. conservativo \Rightarrow se conserva $E = T + U$

$$T = \frac{m \vec{v}^2}{2} = \frac{m R^2 \dot{\varphi}^2}{2} \quad U = -mgR \cos \varphi + \frac{k R^2 (1 + \cos \varphi)^2}{2}$$

$$\frac{m R^2 \dot{\varphi}^2}{2} - mgR \cos \varphi + \frac{k R^2 (1 + \cos \varphi)^2}{2} = E$$

Ec. de mov: $m R^2 \dot{\varphi} \ddot{\varphi} + mgR \operatorname{sen} \varphi \dot{\varphi} + k R^2 (1 + \cos \varphi) (-\operatorname{sen} \varphi \dot{\varphi}) = 0$

$$m R \ddot{\varphi} + mg \operatorname{sen} \varphi + k R (1 + \cos \varphi) \operatorname{sen} \varphi = 0$$

Otra forma: Newton: $m \vec{a} = N \vec{e}_\varphi - B \vec{j} - mg \vec{k} + k R (1 + \cos \varphi) \vec{k}$

$$\vec{a} = R \ddot{\varphi} \vec{e}_\varphi - R \dot{\varphi}^2 \vec{e}_\theta$$

Proyector según \vec{e}_φ : $m R \ddot{\varphi} = -mg \underbrace{\vec{k} \cdot \vec{e}_\varphi}_{\operatorname{sen} \varphi} + k R (1 + \cos \varphi) \underbrace{\vec{k} \cdot \vec{e}_\varphi}_{\operatorname{sen} \varphi}$

Mismo resultado

parte b:

Equilibrio: $\frac{\partial U}{\partial \varphi} = 0$

" $mgR \operatorname{sen} \varphi - \frac{kR^2}{mg} (1 + \cos \varphi) \operatorname{sen} \varphi$

$\frac{\partial U}{\partial \varphi} = mg \operatorname{sen} \varphi (1 - 1 - \cos \varphi) = -mg \operatorname{sen} \varphi \cos \varphi$

1) $\operatorname{sen} \varphi = 0 \Rightarrow \varphi = 0 \text{ o } \pi$

2) $\cos \varphi = 0 \Rightarrow \varphi = \pm \frac{\pi}{2}$

$\frac{\partial^2 U}{\partial \varphi^2} = -mg [\cos^2 \varphi + \operatorname{sen}^2 \varphi]$

1) $\varphi = 0, \pi \quad \frac{\partial^2 U}{\partial \varphi^2} = -mg < 0 \Rightarrow$ Inestable

2) $\varphi = \pm \frac{\pi}{2} \quad \frac{\partial^2 U}{\partial \varphi^2} = mg > 0 \Rightarrow$ Estable

parte d: $\varphi(0) = \frac{\pi}{2}$

$R \dot{\varphi}(0) = v_0 \Rightarrow \dot{\varphi}(0) = \frac{v_0}{R}$

$\frac{mR^2 \dot{\varphi}^2}{2} - mgR \cos \varphi + \frac{kR^2 (1 + \cos \varphi)^2}{2} = \frac{mv_0^2}{2} + \frac{kR^2}{2}$

$-mgR \cos \varphi + \frac{mgR}{2} (1 + 2 \cos \varphi + \cos^2 \varphi)$

$\frac{mR^2 \dot{\varphi}^2}{2} = \frac{mv_0^2}{2} - \frac{mgR}{2} \cos^2 \varphi \geq 0$

$v_0^2 \geq gR \cos^2 \varphi \Rightarrow$ Retroceso en $\cos \varphi_{\text{ret}} = \pm \frac{v_0}{\sqrt{gR}}$

Para que de una vuelta entera $|\cos \varphi_{\text{ret}}| \geq 1$

$|v_0| > \sqrt{gR}$

parte e: $N \geq 0$

Proyecto ec. de Newton sobre \vec{e}_φ

$$-m R \dot{\varphi}^2 = N - mg \underbrace{\vec{k} \cdot \vec{e}_\varphi}_{-\cos\varphi} + kR(1 + \cos\varphi) \underbrace{\vec{k} \cdot \vec{e}_\varphi}_{-\cos\varphi}$$

$$N = mg (1 + \cos\varphi) \cos\varphi - mg \cos\varphi - R \dot{\varphi}^2 m$$

$$N = mg \cos^2\varphi - R \dot{\varphi}^2 m$$

$$\left. \begin{array}{l} \varphi(0) = 0 \\ \dot{\varphi}(0) = \frac{v_1}{R} \end{array} \right\} \Rightarrow N(0) = mg - \frac{v_1^2}{R} m$$

Hay desprendimiento $N < 0 \Rightarrow \boxed{g < \frac{v_1^2}{R}} \Rightarrow \boxed{|v_1| > \sqrt{gR}}$