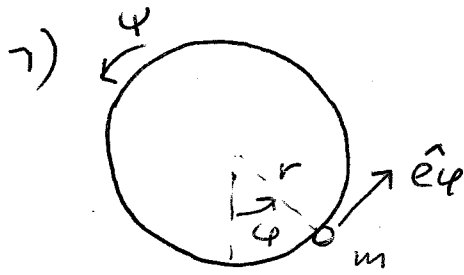


Ej ①



$$\vec{M}_0^{(ext)} \cdot \hat{z} = 0 \quad \forall 0 \in \text{eje}$$

$$\Rightarrow \vec{L}_0 \cdot \hat{z} = L_z = \text{cte.}$$

$$L_z = \frac{1}{2} m r^2 \dot{\varphi} + m r \dot{z} \hat{e}_\varphi$$

$$\vec{v}_m = \dot{z} \hat{z} + r \dot{\varphi} \hat{e}_\varphi = r(\dot{\varphi} + \dot{z}/r) \hat{e}_\varphi + \dot{z} \hat{z}$$

$$\Rightarrow L_z = \frac{1}{2} m r^2 \dot{\varphi} + m r^2 (\dot{\varphi} + \dot{z}/r)$$

$$\underline{L_z = 0 : \ddot{\varphi} + \frac{3}{2} \ddot{z} = 0} \quad (i)$$

Se conserva también la energía:

$$E = mgz + \frac{1}{2} \left(\frac{1}{2} m r^2 \right) \dot{\varphi}^2 + \frac{1}{2} m \left[r^2 (\dot{\varphi} + \dot{z}/r)^2 + \dot{z}^2 \right]$$

$$= mgr\varphi + \frac{1}{2} \left(\frac{1}{2} m r^2 \right) \dot{\varphi}^2 + \frac{1}{2} m r^2 \left[(\dot{\varphi} + \dot{z}/r)^2 + \dot{\varphi}^2 \right]$$

$$\underline{\dot{E} = 0 : \frac{g}{r} \dot{\varphi} + \frac{1}{2} \dot{\varphi} \ddot{\varphi} + \left[(\dot{\varphi} + \dot{z}/r) (\ddot{\varphi} + \ddot{z}/r) + \dot{\varphi} \ddot{\varphi} \right] = 0} \quad (ii)$$

$$\text{Combinando (i) y (ii):} \quad \left\{ \begin{array}{l} \dot{\varphi} = -\frac{3}{4} g/r \\ \ddot{\varphi} = \frac{1}{2} g/r \end{array} \right. \quad (1) \quad (2)$$

$$2) \text{ preintegrando (1) : } \frac{1}{2} (\dot{\varphi}^2 - \dot{\varphi}_0^2) = -\frac{3}{4} g/r \varphi \Leftrightarrow \dot{\varphi}^2 = \dot{\varphi}_0^2 - \frac{3}{2} g/r \varphi$$

¿llega arriba? $z = 2r \Rightarrow \varphi = 2$ si $\dot{\varphi}^2 > 0 \quad \forall \varphi \in [0, 2]$

$$\downarrow$$

$$\underline{\dot{\varphi}_0^2 > 3g/r}$$

3) Integro (1) dos veces en el tiempo:

$$\varphi(t) = \dot{\varphi}_0 t - \frac{3}{8} \frac{g}{r} t^2$$

$$\varphi(t^*) = 2 : t^* = \frac{4}{3} \frac{r}{g} \left(\dot{\varphi}_0 \pm \sqrt{\dot{\varphi}_0^2 - 3g/r} \right)$$

↖ tomo el que me da el menor $t^* > 0$

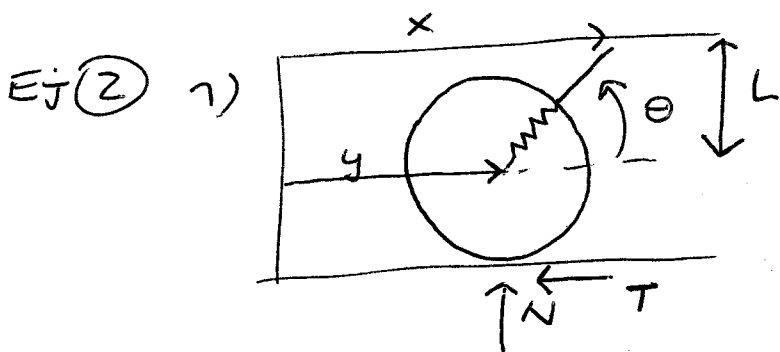
$$\boxed{t^* = \frac{4}{3} \frac{r}{g} \left(\dot{\varphi}_0 - \sqrt{\dot{\varphi}_0^2 - 3g/r} \right)}$$

4) $\vec{v} \cdot \hat{e}_\varphi = 0 : \dot{\varphi}(t^*) + \dot{\psi}(t^*) = 0$

integrando una vez en el tiempo (1) y (2) : $\dot{\varphi} = \dot{\varphi}_0 - 3/4 \frac{g}{r} t$

$$\dot{\psi} = \dot{\psi}_0 + 1/2 \frac{g}{r} t$$

$$\Rightarrow \dot{\varphi}_0 + \dot{\psi}_0 - \frac{7}{4} \frac{g}{r} t^* = 0 : \boxed{\dot{\psi}_0 = -\frac{7}{3} \left(2\dot{\varphi}_0 + \sqrt{\dot{\varphi}_0^2 - 3g/r} \right)}$$



1^{ra} ecuación al disco: $M\ddot{y} = -T + K(x-y)$ (1)

$$0 = N + KL - Mg \quad (2)$$

2^{da} " " : $\frac{1}{2} MR^2 \dot{\omega} = RT$ donde $\dot{\omega} = \dot{y}/R$ (v.s.d.)

$$\frac{1}{2} M \dot{y} = T \quad (3)$$

de (2) : $N = Mg - KL \geq 0$ para que el disco no despegue

sustituyendo (3) en (1) : $2T = -T + K(x-y)$

usando que $\text{tg} \theta = L/(x-y)$

$$T = \frac{KL}{3} \operatorname{ctg} \theta$$

para que ruede sin deslizar: $|T| \leq fN \Leftrightarrow$

$$-f(Mg - KL) \leq \frac{KL}{3} \operatorname{ctg} \theta \leq f(Mg - KL)$$

2) a) Mt. (3) en (1): $M\ddot{y} = -\frac{1}{2}M\ddot{y} + K(x-y)$

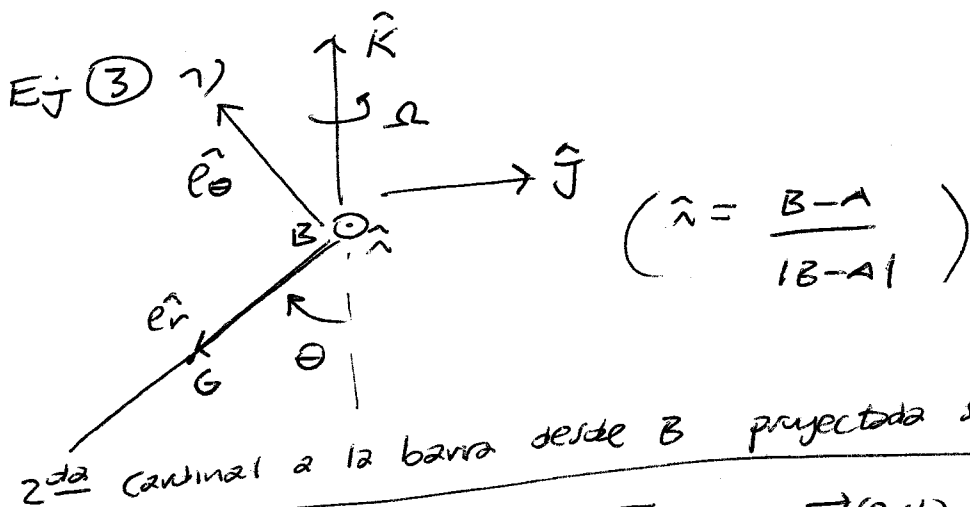
$$\boxed{\frac{3}{2}M\ddot{y} + Ky = Kx(t)}$$

b) $x(t) = x_0 + vt$, la solución gen. a la ec. de movimiento es:

$$y(t) = A \operatorname{sen} \omega t + B \operatorname{cos} \omega t + x_0 + vt, \quad \omega^2 = \frac{2}{3} \frac{K}{M}$$

usando que $y(0) = R$
 $x_0 = R$
 $\dot{y}(0) = 0$

$$\boxed{y(t) = R + v \left(t - \frac{1}{\omega} \operatorname{sen} \omega t \right)}$$



2da. cardinal a la barra desde B proyectada según $\hat{\alpha}$:

$$\left[M(G-B) \times \vec{\omega} + \frac{d}{dt} (\mathbb{I}_B \vec{\omega}) \right] \cdot \hat{\alpha} = \vec{M}_B^{(ext)} \cdot \hat{\alpha} = Mg \frac{L}{2} \operatorname{sen} \theta \quad (1)$$

Consideremos el sistema $\{\hat{\alpha}, \hat{e}_\theta, \hat{e}_r\}$ solidario al rígido y principal

$$\frac{d}{dt} (\mathbb{I}_B \vec{\omega}) = \mathbb{I}_B \dot{\vec{\omega}} + \vec{\omega} \times (\mathbb{I}_B \vec{\omega})$$

$$\text{siendo } \vec{\omega} = -\dot{\theta} \hat{\alpha} + \Omega \hat{K} = -\dot{\theta} \hat{\alpha} + \Omega \operatorname{sen} \theta \hat{e}_\theta - \Omega \operatorname{cos} \theta \hat{e}_r$$

$$\vec{\omega} = -\ddot{\theta} \hat{\lambda} + (\dot{\Omega} \sin\theta + \Omega \cos\theta \dot{\theta}) \hat{e}_\theta - (\dot{\Omega} \cos\theta - \Omega \sin\theta \dot{\theta}) \hat{e}_r$$

$$I_B \{ \hat{\lambda}, \hat{e}_\theta, \hat{e}_r \} = \begin{pmatrix} I & 0 & 0 \\ 0 & I & 0 \\ 0 & 0 & 0 \end{pmatrix}, \quad I = \int_0^L \left(\frac{M}{L}\right) x^2 dx = \frac{1}{3} ML^2$$

$$\Pi_B \vec{\omega} \cdot \hat{\lambda} = -I \ddot{\theta}$$

$$\Pi_B \vec{\omega} = I(-\ddot{\theta} \hat{\lambda} + \Omega \sin\theta \hat{e}_\theta), \quad \vec{\omega} \times \Pi_B \vec{\omega} \cdot \hat{\lambda} = I \Omega^2 \sin\theta \cos\theta$$

$$\vec{\alpha}_B = L \dot{\Omega} \hat{j} - L \Omega^2 \hat{\lambda} = L \dot{\Omega} (-\cos\theta \hat{e}_\theta - \sin\theta \hat{e}_r) - L \Omega^2 \hat{\lambda} \quad \Rightarrow$$

$$G - B = \frac{L}{2} \hat{e}_r$$

$$\Rightarrow M(G - B) \times \vec{\alpha}_B \cdot \hat{\lambda} = \frac{ML^2}{2} \dot{\Omega} \cos\theta$$

y la ec. de movimiento (7) es:

$$\ddot{\theta} + \sin\theta \left(\frac{3}{2} g/L - \Omega^2 \cos\theta \right) - \frac{3}{2} \dot{\Omega} \cos\theta = 0$$

$$2) \Omega = \Omega_0 : \ddot{\theta} + \sin\theta \left(\frac{3}{2} g/L - \Omega_0^2 \cos\theta \right) = 0$$

$$\text{equilibrio } (\ddot{\theta} = 0) : \sin\theta = 0 \quad \theta = 0, \pi$$

$$\cos\theta = \frac{3}{2} g/L \Omega_0^{-2} \quad \exists \text{ si } \Omega_0^2 > \frac{3}{2} g/L$$

Los ptos. de eq. estable son los mínimos de U' / $\frac{dU'}{d\theta} = \sin\theta \left(\frac{3}{2} g/L - \Omega_0^2 \cos\theta \right)$

$$\frac{d^2U'}{d\theta^2} = \cos\theta \left(\frac{3}{2} g/L - \Omega_0^2 \cos\theta \right) + \Omega_0^2 \sin^2\theta$$

$$\left. \frac{d^2U'}{d\theta^2} \right|_{\pi} = - \left(\frac{3}{2} g/L + \Omega_0^2 \right) < 0 \Rightarrow \theta = \pi \text{ inestable}$$

$$\left. \frac{d^2U'}{d\theta^2} \right|_0 = \frac{3}{2} g/L - \Omega_0^2 > 0 \text{ si } \Omega_0^2 < \frac{3}{2} g/L$$

$$\frac{d^2U^1}{d\theta^2} \Big|_{\cos\theta = \frac{3}{2}g/L \cdot \Omega_0^{-2}} = \Omega_0^2 \left(1 - \left(\frac{3}{2}g/L \cdot \Omega_0^{-2} \right)^2 \right) > 0 \text{ si}$$

$$\Omega_0^2 > \frac{3}{2}g/L$$

Para $\Omega_0^2 = \frac{3}{2}g/L$: $\frac{dU^1}{d\theta} = \Omega_0^2 \sin\theta (1 - \cos\theta)$

$\theta \rightarrow 0^+ \quad \frac{dU^1}{d\theta} > 0$

$\theta \rightarrow 0^- \quad \frac{dU^1}{d\theta} < 0$

$\Rightarrow \frac{dU^1}{d\theta}$ creciente alrededor de $\theta=0$: $\theta=0$ es un mínimo

resumen:

$\theta=0$ estable si $\Omega_0^2 \leq \frac{3}{2}g/L$

$\theta=\pi$ inestable

$\cos\theta = \frac{3}{2}g/L \cdot \Omega_0^{-2} \exists$ y es estable para $\Omega_0^2 > \frac{3}{2}g/L$